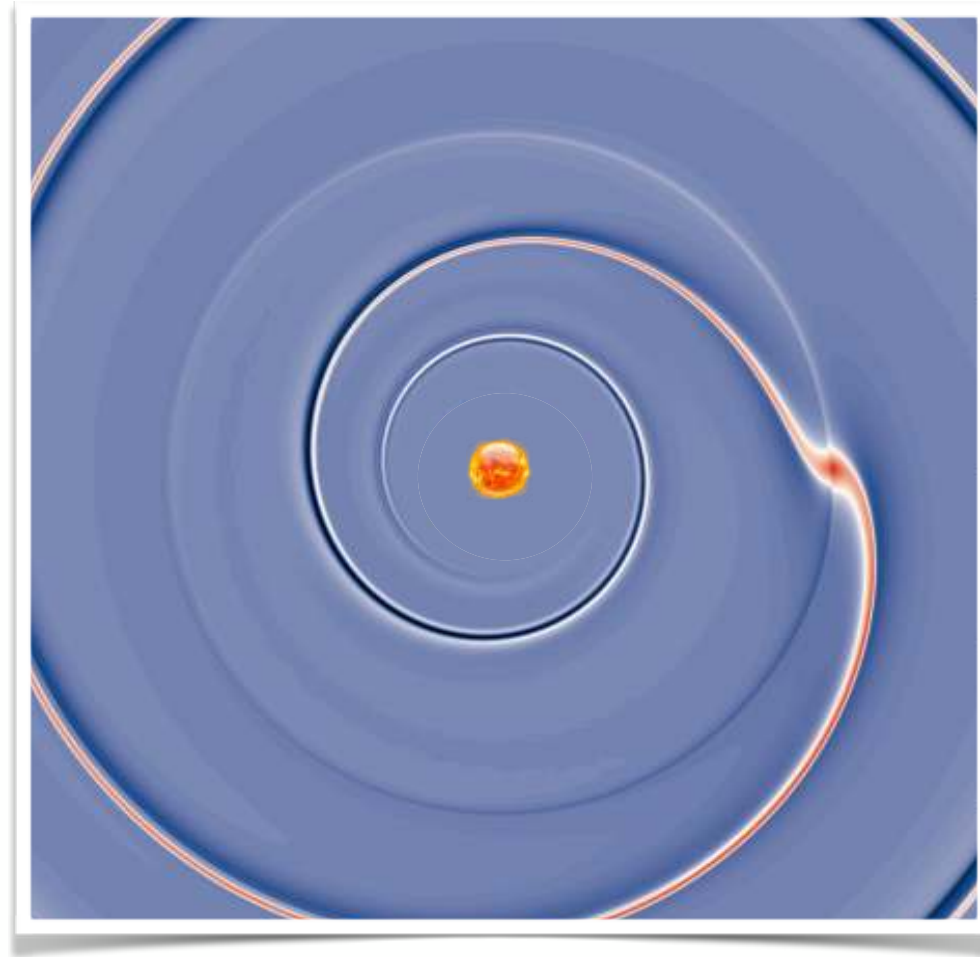


# Lecture 4: disk-planet interactions and planetary migration



## *Suggested references:*

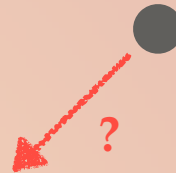
- Baruteau et al. 2016, **Formation, Orbital and Internal Evolutions of Young Planetary Systems**  
[arxiv.org/abs/1604.07558](https://arxiv.org/abs/1604.07558)
- Baruteau et al. 2014, **Planet-Disc Interactions and Early Evolution of Planetary Systems**  
[arxiv.org/abs/1312.4293](https://arxiv.org/abs/1312.4293)

# Planet formation and orbital evolution

## planet-disc interactions

change planets semi-major axes  
(planetary *migration*)

damp eccentricities and inclinations

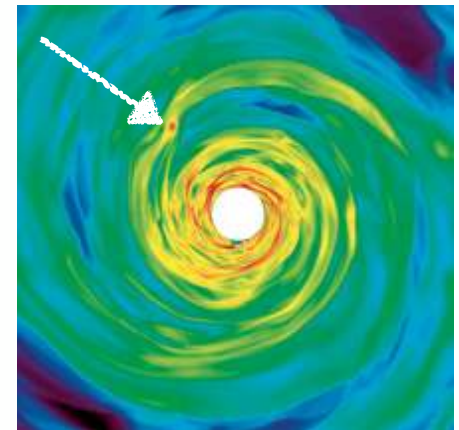


## planet formation

core accretion?



gravitational instability?



protoplanetary disc

# Planet formation and orbital evolution

## planet-disc interactions

change planets semi-major axes  
(planetary *migration*)

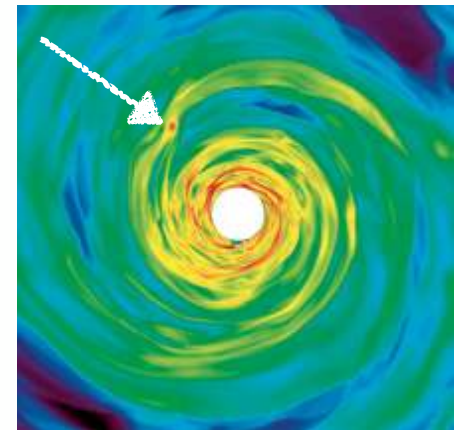
damp eccentricities and inclinations

## planet formation

core accretion?



gravitational instability?



## planet-planet interactions

also change semi-major axes!  
pump eccentricities and inclinations  
protoplanetary disc

# Planet formation and orbital evolution

## disc dispersal

(after 1-10 Myr)

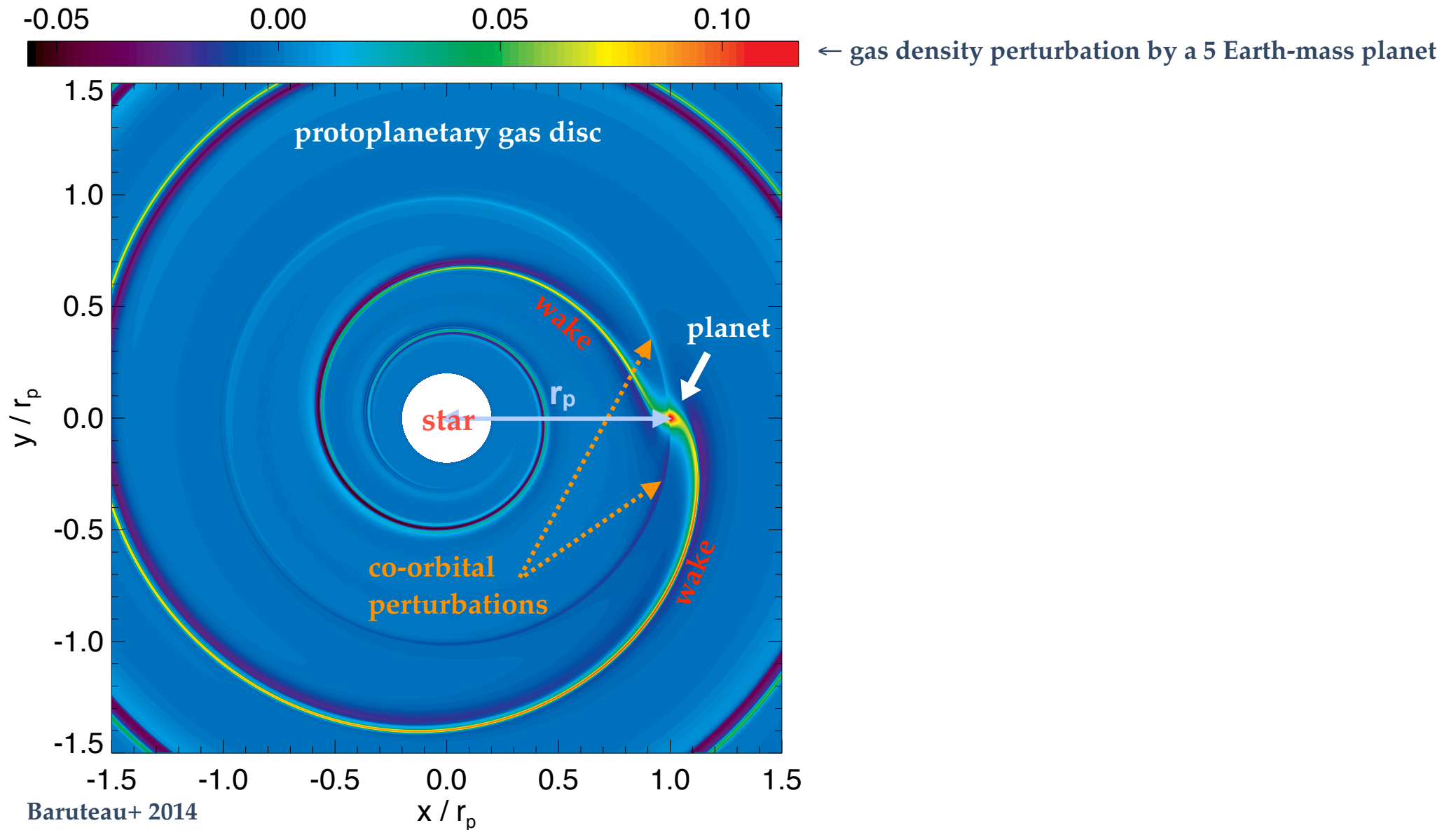
- . interactions with the central star (tides, stellar evolution) or with nearby stars
- . planet-planet interactions
- . planets-debris disc interactions





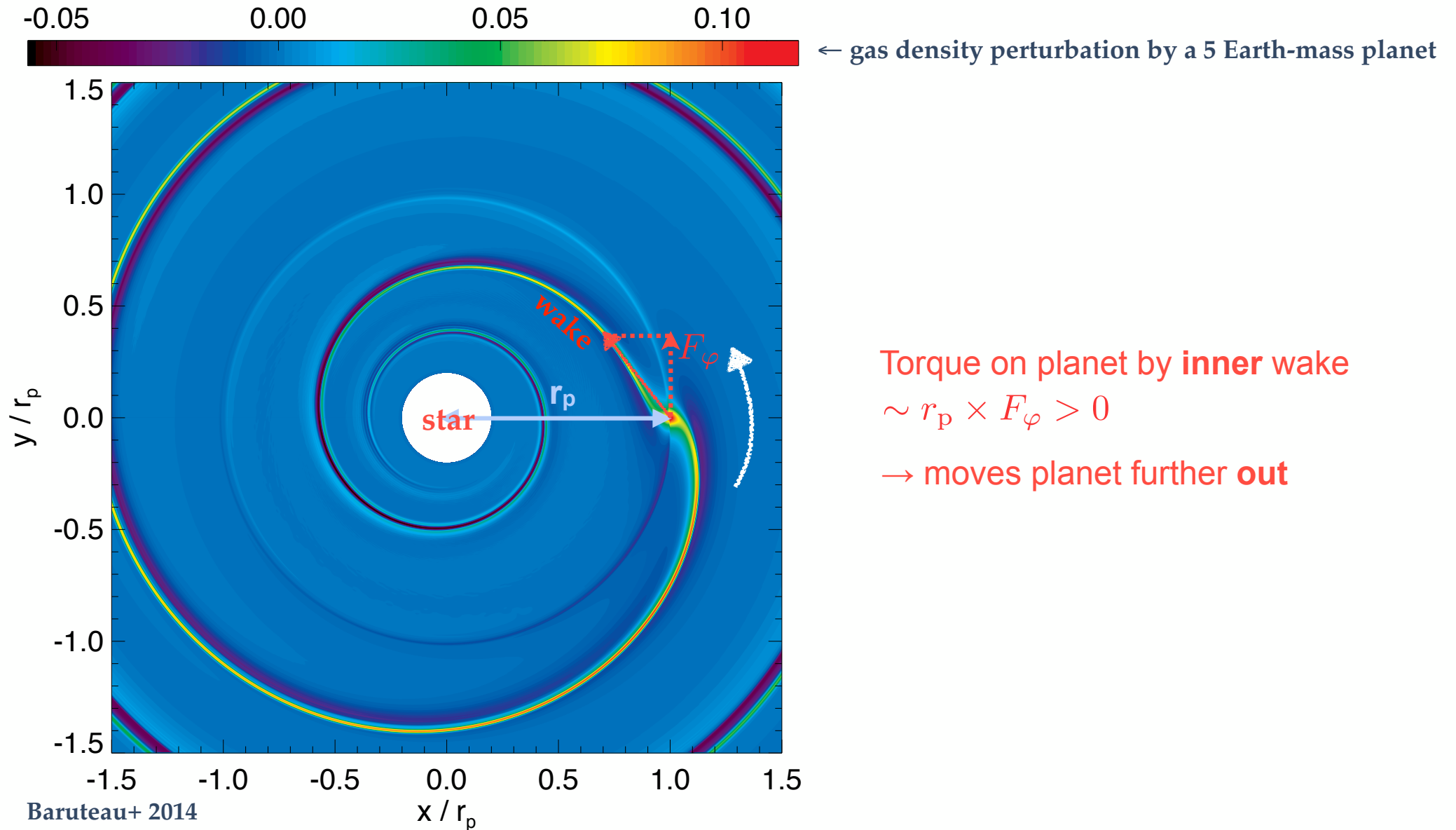
# Disc migration of low-mass planets

(typically up to  $\sim 10$  Earth masses)



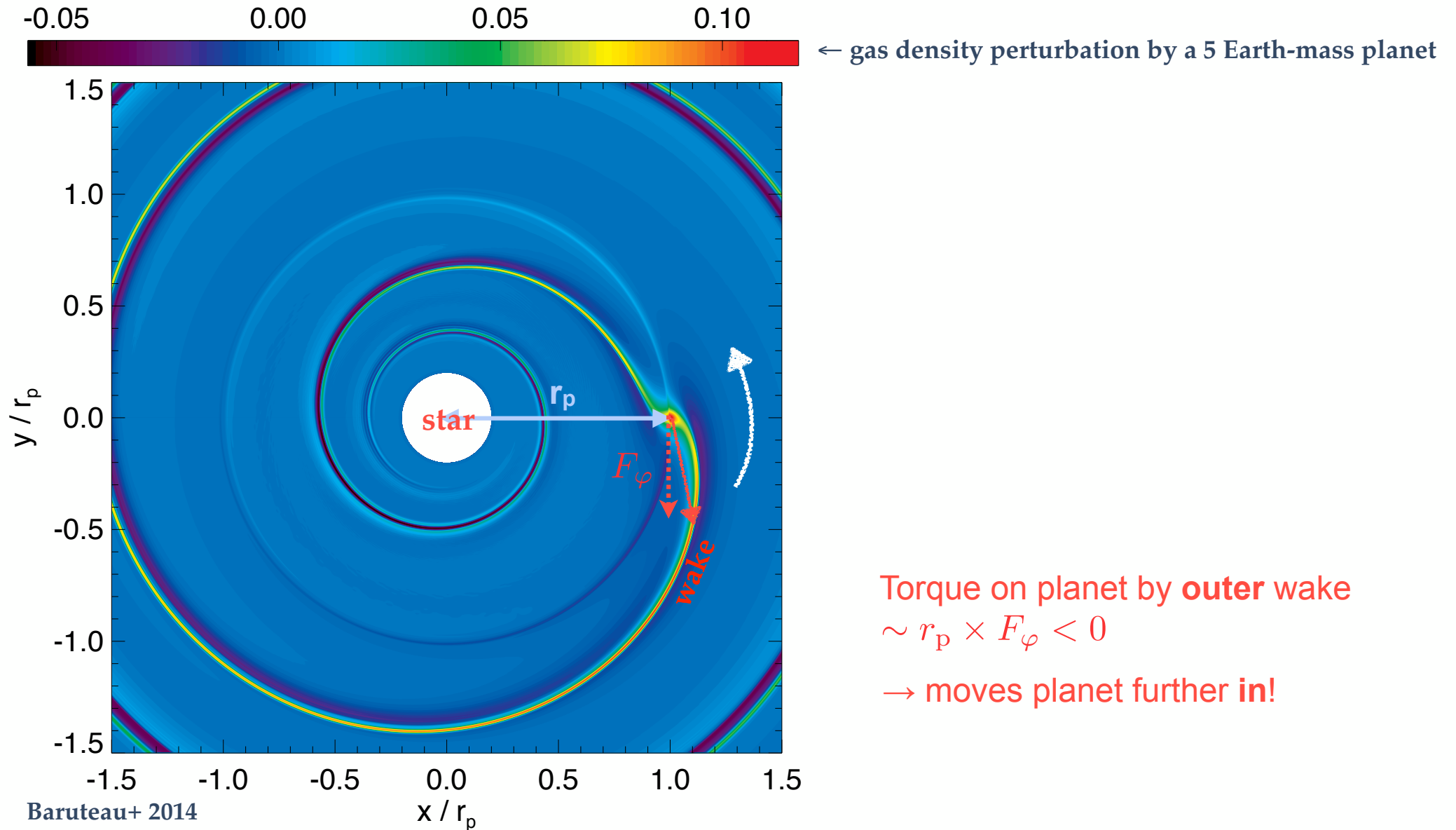
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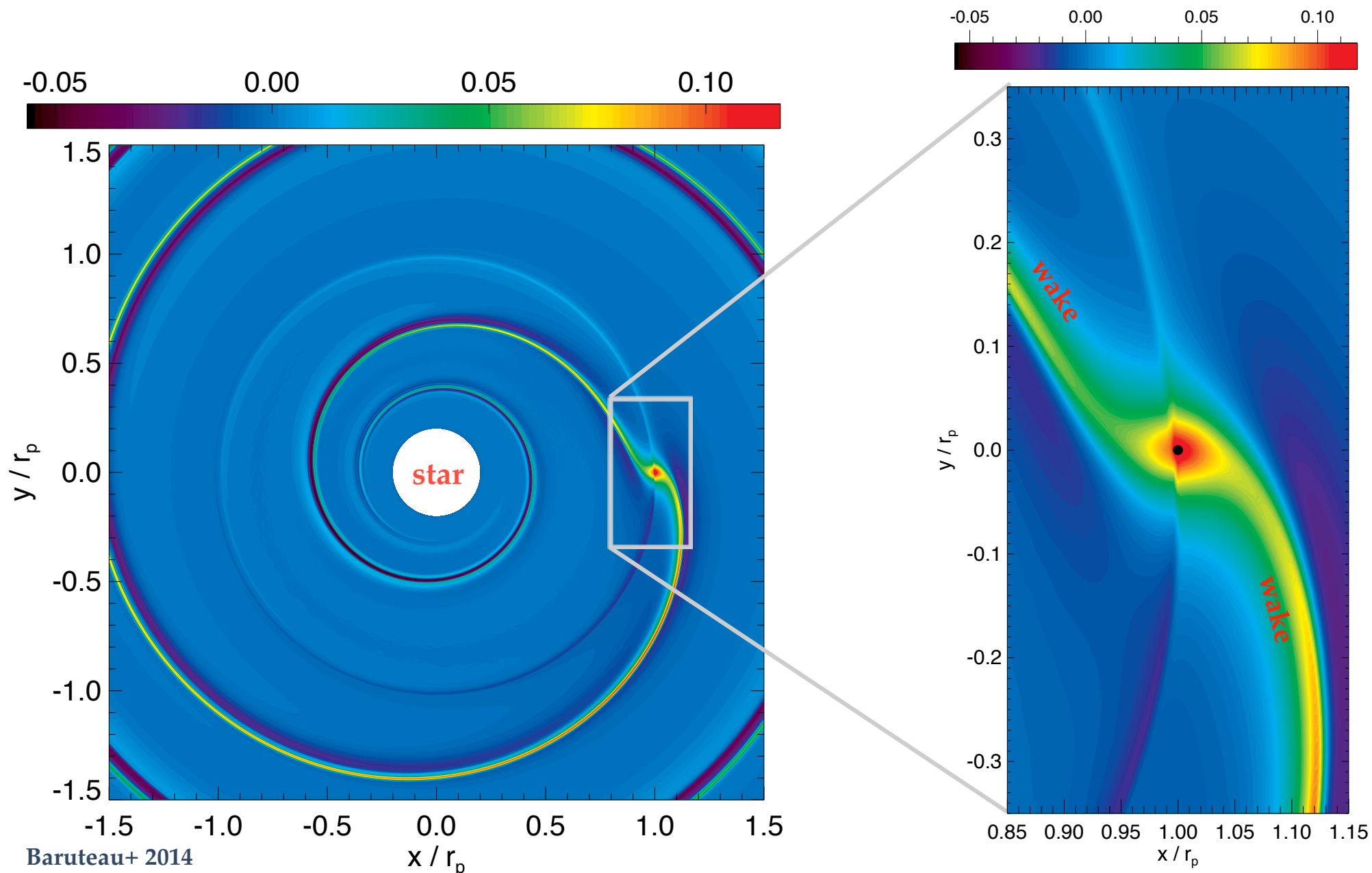
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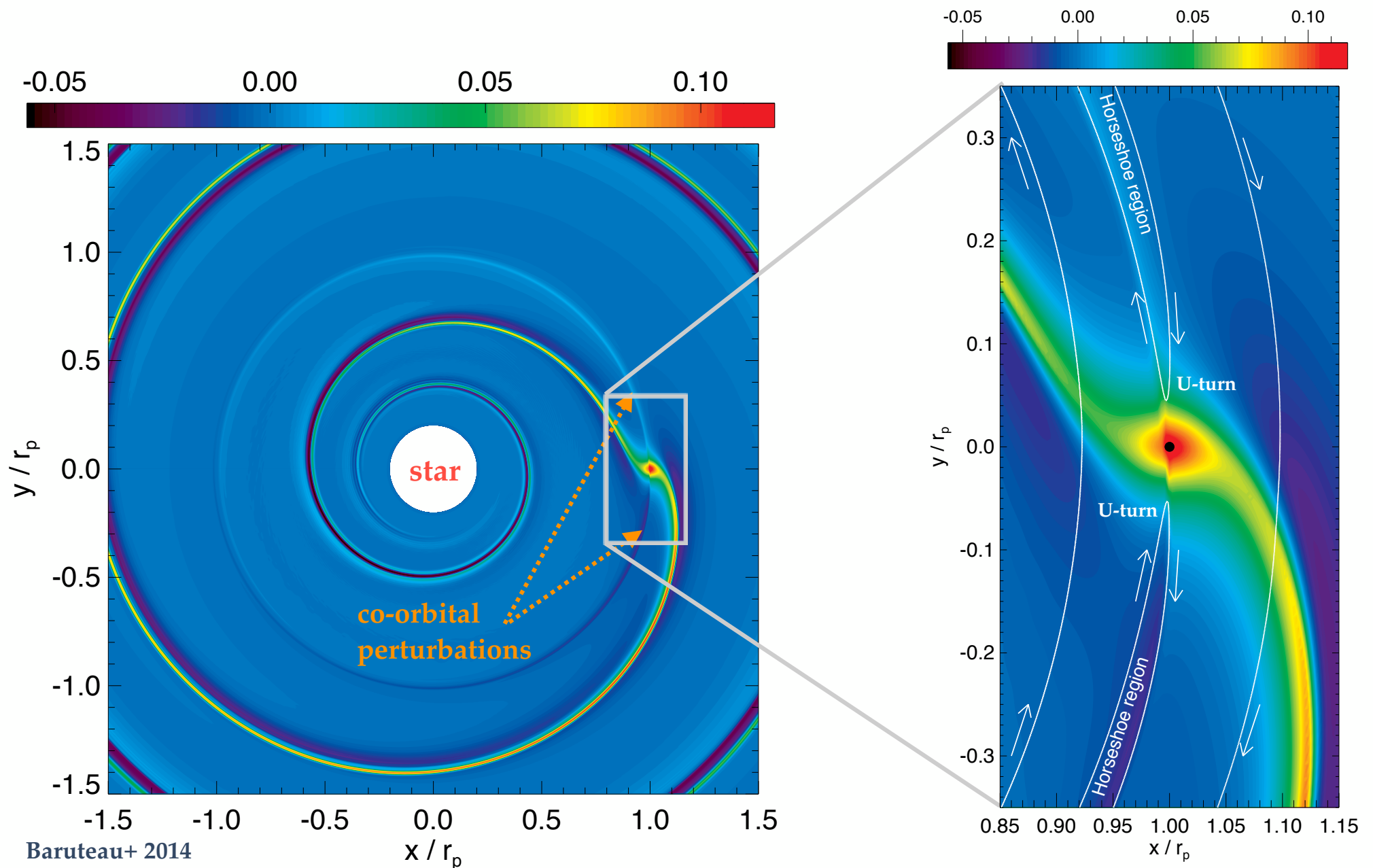
Baruteau+ 2014

The total wake torque is (generally) negative and favors inward migration



# Disc migration of low-mass planets

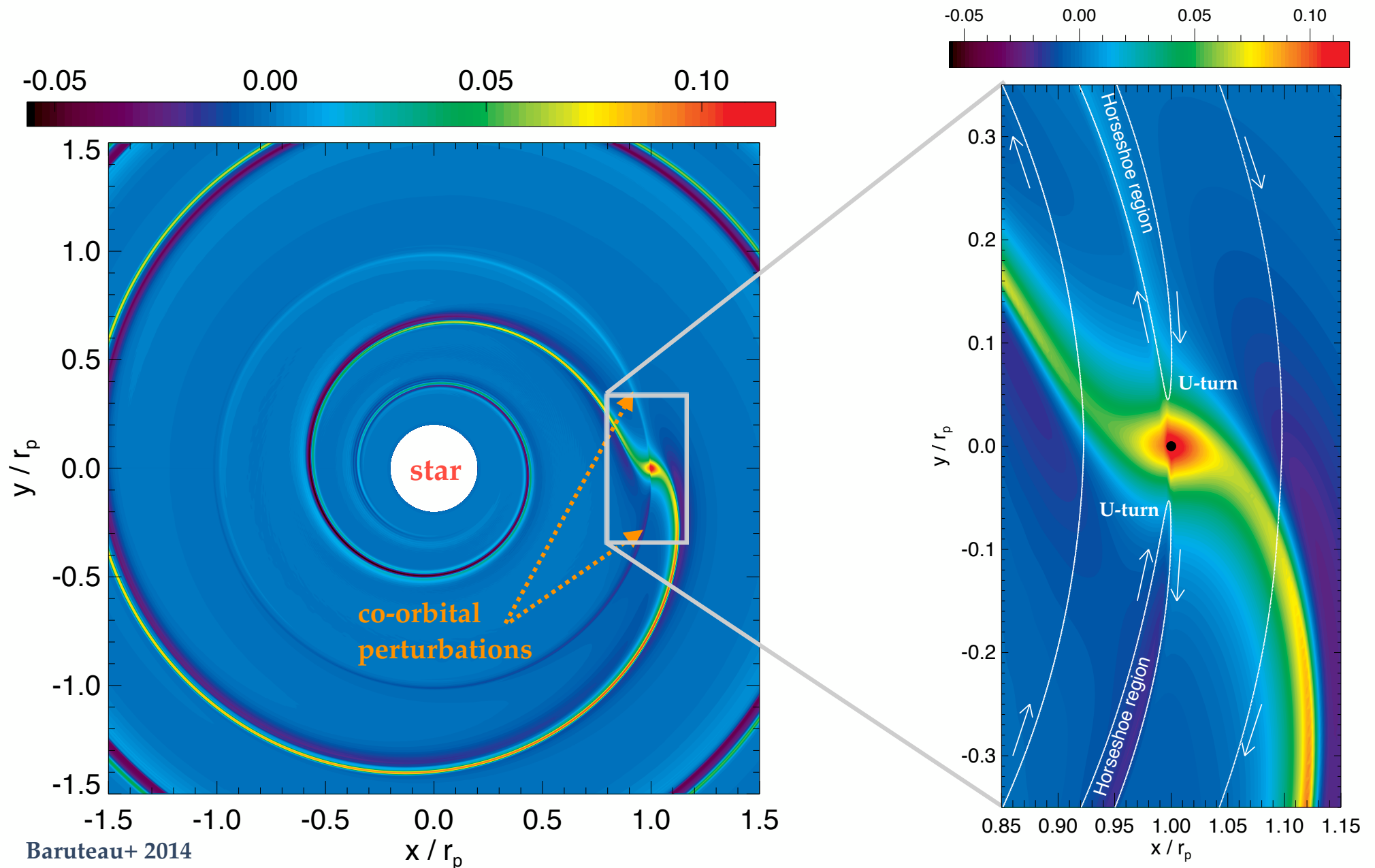
(typically up to  $\sim 10$  Earth masses)



The horseshoe region exerts a torque on the planet called **corotation torque**

# Disc migration of low-mass planets

(typically up to  $\sim 10$  Earth masses)



The corotation torque is generally **positive** and favors **outward migration**



# Focus on the planet wakes

- They are the superposition of spiral density waves emitted at **Lindblad resonances**
- linear problem: **Lindblad resonances** = where the gas azimuthal velocity *relative to the planet* matches  $\pm$  the phase velocity of acoustic waves in the azimuthal direction

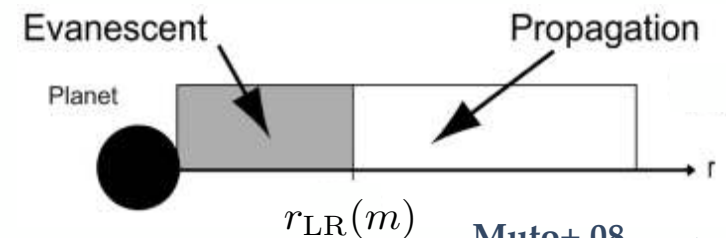
$$\begin{aligned} \text{Dopper-shifted wave frequency } \omega & \rightarrow \Omega - \Omega_p = \pm \frac{\omega}{m} \\ \text{azimuthal wavevector} = m/R & \rightarrow \omega = \pm m(\Omega - \Omega_p) \\ \text{horizontal epicyclic frequency } \kappa \approx \Omega & \rightarrow \omega^2 = \kappa^2 + m^2 c_s^2 / R^2 \end{aligned} \Rightarrow r_{\text{LR}} = r_p \left( 1 \pm \sqrt{h^2 + m^{-2}} \right)^{2/3}$$

disk's aspect ratio

- ▶ the 2/3 power implies that **outer** resonances lie slightly **closer** to the planet than **inner** resonances
- ▶ at large  $m$ ,  $v_{\text{phase}} \rightarrow c_s$  and Lindblad resonances pile up at  $\pm 2H/3$  around the planet
- ▶ 2D + **WKB** approximation, the wave equation reduces to a Schrödinger-like equation **for every  $m$** :

$$\frac{d^2 f}{dr^2} + V(r)f(r) = S(r)$$

$f(r)$   $\equiv$  gas perturbed enthalpy,  
 $S(r)$  = forcing term due to planet's potential  
 $V(r)$  defines regions of wave propagation:  
 $V(r) > 0$ : waves can **propagate**  
 $V(r) < 0$ : waves are **evanescent**  
 $V(r) = 0$  defines **resonances** location



# Focus on the planet wakes

- They are the superposition of spiral density waves emitted at **Lindblad** resonances
- ❖ linear problem: **Lindblad resonances** = where the gas azimuthal velocity *relative to the planet* matches  $\pm$  the phase velocity of acoustic waves in the azimuthal direction

$$\begin{aligned} \triangleright v_\varphi - v_p = \pm \frac{\omega}{k_\varphi} &\rightarrow \Omega - \Omega_p = \pm \frac{\omega}{m} \\ \text{Dopper-shifted wave frequency } \omega & \\ \text{azimuthal wavevector } = m/R & \end{aligned} \quad \left. \begin{array}{l} \text{For a gas disk without self-gravity, } \omega^2 = \kappa^2 + m^2 c_s^2 / R^2 \\ \text{horizontal epicyclic frequency } \kappa \approx \Omega \end{array} \right\} \Rightarrow r_{\text{LR}} = r_p \left( 1 \pm \sqrt{h^2 + m^{-2}} \right)^{2/3}$$

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- ▶ the location of Lindblad resonances depends on whatever changes  $\Omega$ ,  $\Omega_p$  or the **waves dispersion relation**, like the disk's self-gravity, magnetic field...

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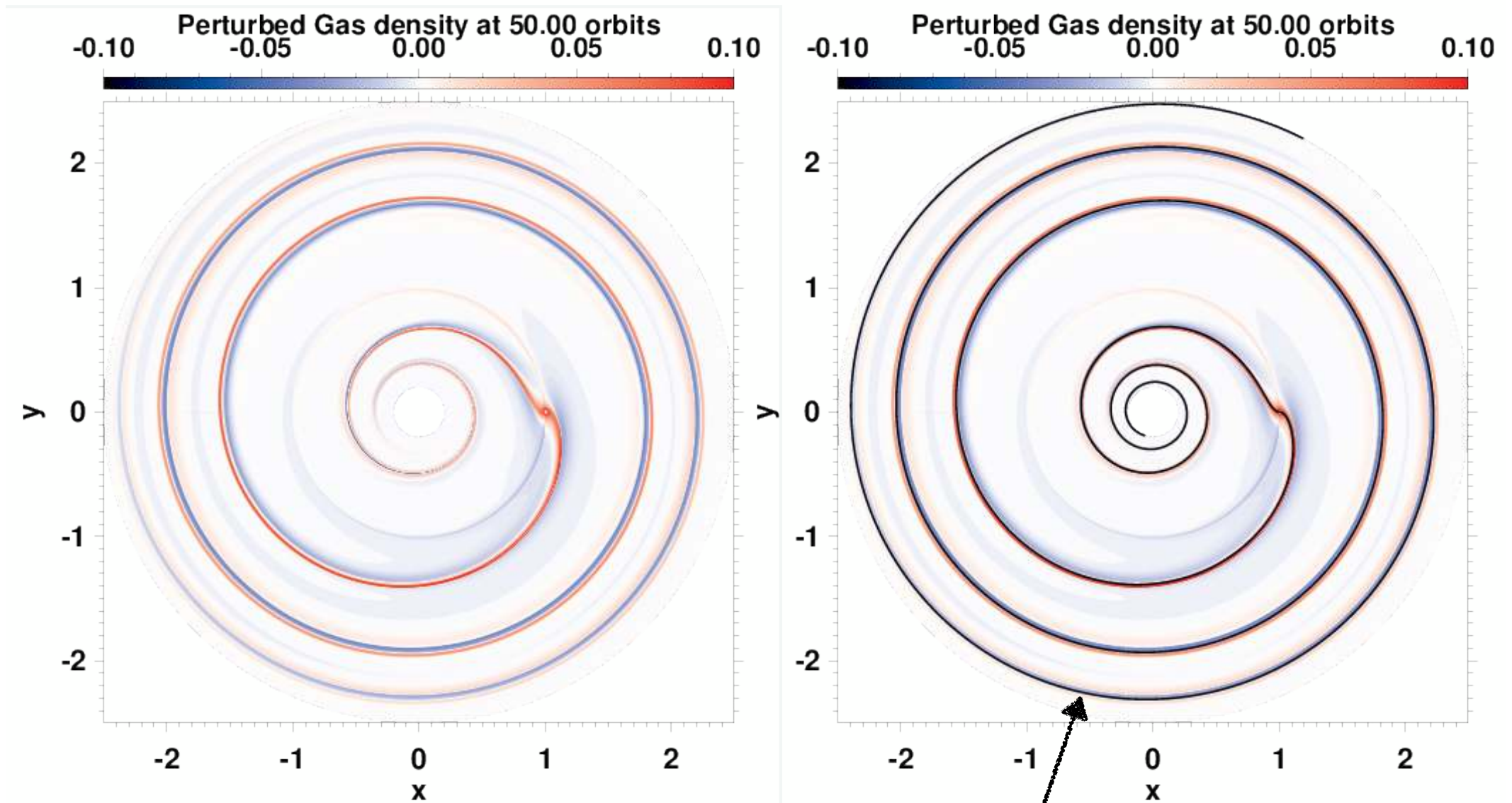
For a gas disk without self-gravity,  $\omega^2 = \kappa^2 + m^2 c_s^2 / R^2$

horizontal epicyclic frequency  $\kappa \approx \Omega$  (arrow to  $\kappa^2$ )

disk's aspect ratio (arrow to  $h$ )

- ▶ the 2/3 power implies that **outer** resonances lie slightly **closer** to the planet than **inner** resonances
- ▶ at large  $m$ ,  $v_{\text{phase}} \rightarrow c_s$  and Lindblad resonances pile up at  $\pm 2H/3$  around the planet
- ▶ 2D + **WKB** approximation, the wave equation reduces to a Schrödinger-like equation **for every  $m$** :
- ▶ the location of Lindblad resonances depends on whatever changes  $\Omega$ ,  $\Omega_p$  or the **waves dispersion relation**, like the disk's self-gravity, magnetic field...
- ▶ waves launched at Lindblad resonances **interfere constructively** into an **one-armed spiral wave**, called the planet's **wake**, which co-rotates with the planet

# Focus on the planet wakes



waveform based on linear theory  
(Ogilvie & Lubow 02)



Thanks for your attention!

